# Simultaneous measurement of intra- and intermolecular NOEs in differentially labeled protein-ligand complexes* 

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#### Abstract

A new NOE strategy is presented that allows the simultaneous observation of intermolecular and intramolecular NOEs between an unlabeled ligand and a ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled protein. The method uses an adiabatic ${ }^{13} \mathrm{C}$ inversion pulse optimized to an empirically observed relationship between ${ }^{1} J_{\mathrm{CH}}$ and carbon chemical shift to selectively invert the protein protons (attached to ${ }^{13} \mathrm{C}$ ). Two NOESY data sets are recorded where the intermolecular and intramolecular NOESY cross peaks have either equal or opposite signs, respectively. Addition and subtraction yield two NOESY spectra which contain either NOEs within the labeled protein (or unlabeled ligand) or along the binding interface. The method is demonstrated with an application to the $\mathrm{B}_{12}$-binding subunit of Glutamate Mutase from Clostridium tetanomorphum complexed with the $\mathrm{B}_{12}$-nucleotide loop moiety of the natural cofactor adenosylcobalamin (Coenzyme $\mathrm{B}_{12}$ ).


## Introduction

NMR spectroscopy has become a powerful tool to study biologically important noncovalent assemblies of macromolecules in solution and yielded a wealth of information on the structural details of molecular recognition (Clore and Gronenborn, 1997). The most efficient experimental strategy relies on uniform labeling of one of the components with ${ }^{15} \mathrm{~N}$ and/or ${ }^{13} \mathrm{C}$, while leaving the second component unlabeled (Otting et al., 1986; Otting and Wüthrich, 1990; Sattler et al., 1999; Breeze, 2000). This approach allows the selection of intermolecular NOEs between the two constituents of the macromolecular complex (Wider et al., 1990, 1991; Folkers et al., 1993, 1994; Bax et al., 1994; Yang et al., 1995). The accuracy and precision of the complex structure is particularly sensitive to the number of intermolecular NOE connectivities, thus making the NOE assign-

[^0]ment the essential (and time consuming) part of the structure elucidation process (Clore, 2000). Specifically, the editing/filtering NMR experiments devised for this purpose employ the one-bond heteronuclear scalar couplings to select for/against protons that are bound at sites carrying an isotope label. A major hurdle to a general application of the method is due to variations in JCH scalar coupling constants (between $120-220 \mathrm{~Hz}$ ). This problem has been overcome (in part) by developing second-(high)-order isotope filters (Gemmecker et al., 1992; Ikura and Bax, 1992; Lee et al., 1994; Dalvit et al., 1998; 1999; Sattler et al., 1999; Breeze, 2000), in particular by means of adiabatic inversion pulses with a sweep rate adjusted to an empirical $J_{\mathrm{CH}}$ vs ${ }^{13} \mathrm{C}$ chemical shift profile (Zwahlen et al., 1997) or based on a $J$-compensated isotope filter through isomorphism to composite pulse rotations (Stuart et al., 1999). The drawback of loosing all intramolecular NOEs of the labeled macromolecule was recently overcome by a $J$-spectroscopy based 4D experiment for simultaneous recording of inter- and intramolecular NOEs (Melacini, 2000).

Here we report a filtering/editing technique based on selective inversion of protons bound to ${ }^{13} \mathrm{C}$ (using adiabatic pulses with optimized sweep rates) and ${ }^{15} \mathrm{~N}$, and with subsequent addition and subtraction of the 3D NOESY data sets to simultaneously identify intraand intermolecular NOE connectivities. The filtering strategy is related to the 'linear-combination'-type filters filters and edits introduced by Wüthrich and coworkers (Wörgötter et al., 1986; Otting et al., 1986; Otting and Wüthrich, 1990). The proposed sequence provides good NOE selectivity and offers an improvement in sensitivity for the measurement of intramolecular NOEs relative to the previously devised semiconstant J-resolved 4D experiment (Melacini, 2000). The method is applied to the $\mathrm{B}_{12}$-binding subunit of Glutamate Mutase from Clostridium tetanomorphum complexed with the $\mathrm{B}_{12}$-nucleotide loop moiety of the natural cofactor Coenzyme $\mathrm{B}_{12}$.

## Materials and methods

All NMR experiments were performed on a Varian UNITYPlus 500 MHz spectrometer equipped with a pulse field gradient unit and triple resonance probes with actively shielded $z$ gradients. All spectra were recorded at $26{ }^{\circ} \mathrm{C} .{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled MutS was overexpressed in E. coli and purified as described previously (Hoffmann et al., 2001). The NMR sample contained $2.0 \mathrm{mM}{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled MutS, 5 mM DTT, 1 mM EDTA, $0.03 \% \mathrm{NaN}_{3}$, and 10 mM phosphate buffer at pH 6.9, dissolved in $90 \% \mathrm{H}_{2} \mathrm{O} / 10 \% \mathrm{D}_{2} \mathrm{O}$. In the binary complex, the ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ MutS: $\mathrm{B}_{12}$-nucleotide ratio was 1.0:12.0. The $\mathrm{B}_{12}$-nucleotide was prepared as described elsewhere (Eschenmoser, 1988; Bartels et al., 1991). The dissociation constant of the complex $K_{d}$ was determined from a titration of the protein with the $\mathrm{B}_{12}$-nucleotide by a non-linear least squares fitting of the chemical shift changes as a function of ligand concentration ( $K_{d}=5.6 \pm 0.7 \mathrm{mM}$; Tollinger et al., 2001). The 3D SIM-NOESY-HSQC was recorded with the pulse scheme of Figure 3 ( 180 ms NOESY mixing time). Intra- and intermolecular NOESY spectra were obtained by post-acquisition addition and subtraction of the data (recorded in an interleaved manner). A data matrix consisting of 45,35 and 512 complex points in each of $t_{1}, t_{2}$ and $t_{3}$ was acquired, the following spectral windows were used: F1: 6000 Hz , F2: 3000 Hz and F3: 8000 Hz . A relaxation delay of 1 s was used along with 16 scans per FID. For comparison a $150 \mathrm{~ms} \mathrm{3D}{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ NOESY-HSQC (Pascal


Figure 1. J-coupling evolution scheme for selective inversion of protons that are scalar coupled to ${ }^{15} \mathrm{~N}$ or ${ }^{13} \mathrm{C}$ spins. The rf carrier frequency of the ${ }^{13} \mathrm{C}$ inversion pulse is modified. Carbon magnetizations of carbons resonating at $\omega_{\mathrm{a}}$ and $\omega_{\mathrm{b}}$ are inverted at different times $t_{a}$ and $t_{b}$. The rate of the frequency sweep $(\mathrm{d} \omega / \mathrm{d} t)$ is optimized with respect to the ${ }^{1} J_{\mathrm{CH}}$ vs $\delta_{\mathrm{C}}$ relationship. $2 \tau_{\mathrm{c}}=1 / J_{\mathrm{NH}}$, so that protons bound to ${ }^{15} \mathrm{~N}$ will be inverted. The net evolution time during which the heteronuclear $J_{\mathrm{CH}}$ scalar coupling evolves is $2 \tau_{c}-2 \tau_{\mathrm{a}}$ (for carbon spin with resonance frequency $\omega_{\mathrm{a}}$ ) or $2 \tau_{\mathrm{c}}-2 \tau_{\mathrm{b}}$ (for carbon spin with resonance frequency $\omega_{\mathrm{b}}$ ).
et al., 1994) was recorded. A data matrix consisting of 45,50 and 512 complex points in each of $t_{1}$, $t_{2}$ and $t_{3}$ was acquired, the following spectral windows were used: F1: $6000 \mathrm{~Hz}, \mathrm{~F} 2: 3000 \mathrm{~Hz}$ and F3: 8000 Hz . Again, a relaxation delay of 1 s and 16 scans were used. All spectra were processed identically using NMRPipe/NMRDraw software (Delaglio et al., 1995) and analyzed using the program NMRView (Johnson and Blevins, 1994). The sizes of the two 3D HSQC-NOESY data set were doubled in both indirect frequency dimensions and the data set were apodized in each dimensions with a squared sinebell window function, zero filled, Fourier transformed, phased, and the imaginaries eliminated. The final data set comprised $256 \times 128 \times 512$ real points.

Figure 1 illustrates the basic element which is used to selectively invert the protons bound to ${ }^{13} \mathrm{C}$ or ${ }^{15} \mathrm{~N}$. The overall delay $2 \tau_{\mathrm{c}}$ is set to $2 \tau_{\mathrm{c}}=1 / J_{\mathrm{NH}}$, so that protons bound to ${ }^{15} \mathrm{~N}$ will be inverted. Due to the dependence of $J_{\mathrm{CH}}$ on ${ }^{13} \mathrm{C}$ chemical shift a different strategy is employed for inversion of protons bound to ${ }^{13} \mathrm{C}$. The effect of an adiabatic inversion pulse can be appreciated by considering that spin inversion occurs when the frequency sweep passes through the resonance frequency of the carbon of interest, provided that the overall sweep width is large compared to the radio-frequency (rf) field. We consider two carbons with their resonance frequencies $\omega_{\mathrm{a}}$ and $\omega_{\mathrm{b}}$, respectively. If we start the frequency sweep from the
up-field side of the spectrum, the accelerating reference (spin-lock) frame will pass through the resonance frequency $\omega_{\mathrm{a}}$ at time point $t_{a}$, whereas the frequency $\omega_{\mathrm{b}}$ will be matched at a later time $t_{b}$. Thus the time when the two carbons are inverted will be different, resulting in a different net evolution time during which the heteronuclear JCH scalar coupling evolves $\left(2 \tau_{\mathrm{c}}-2 \tau_{\mathrm{a}}\right.$ or $\left.2 \tau_{\mathrm{c}}-2 \tau_{\mathrm{b}}\right)$. A substantial improvement in ${ }^{1} \mathrm{H}$ inversion efficiency can be achieved by adjusting the sweep rate of the adiabatic pulse to the almost linear relation between the one bond ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ scalar coupling, $J_{\mathrm{CH}}$, and carbon chemical shift, $\delta_{\mathrm{C}},\left(J_{\mathrm{CH}}=\right.$ $A \delta_{\mathrm{C}}+B$, where $A=0.71 \mathrm{~Hz} / \mathrm{ppm} ; B=101 \mathrm{~Hz}$ ) (Zwahlen et al., 1997). Figure 2B shows the improvement obtained by $J_{\mathrm{CH}}$ compensation using the scheme of Figure 1 compared to a conventional INEPT-type echo sequence (Figure 2A) with $2 \tau_{c}=1 / J_{\mathrm{CH}}$. Although the linear $J_{\mathrm{CH}}$ vs ${ }^{13} \mathrm{C}$ chemical shift profile can be best accomplished by a hyperbolic $\omega_{\mathrm{rf}}(t)$ profile, we used optimized linear frequency sweeps as described elsewhere (Zwahlen et al., 1997). The outcome of the basic sequence element of Figure 1 was determined from density matrix calculations assuming a heteronuclear two-spin system. In all calculations a total sweep width of 60 kHz and adiabatic pulse durations of $2.8 \mathrm{~ms}\left(\gamma \mathrm{~B}_{1}: 5.0 \mathrm{kHz}\right)$ were employed.

The 3D SIM-NOESY-HSQC experiment for simultaneous measurement of inter- and intramolecular NOEs is outlined in Figure 3. In terms of the flow of magnetization it is analogous to the ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-edited NOESY-HSQC experiment of Kay and co-workers (Pascal et al., 1994), however, preceded by the pulse sequence element of Figure 1. Two data sets are recorded where in one experiment the heteronuclear ${ }^{13} \mathrm{C}$ and ${ }^{15} \mathrm{~N}$ inversion pulses are omitted and the onebond ${ }^{1} J_{\mathrm{CH}}$ and ${ }^{1} J_{\mathrm{NH}}$ scalar couplings are refocused. The $90^{\circ}{ }^{1} \mathrm{H}$ radio-frequency (rf) pulse (before the NOESY period $\tau_{m}$ ) returns both ligand and protein ${ }^{1} \mathrm{H}$ transverse magnetizations to the $z$-axis. In contrast, in the second experiment the pulse scheme of Figure 1 is applied and both ${ }^{1} J_{\mathrm{CH}}$ and ${ }^{1} J_{\mathrm{NH}}$ couplings are active during the time $2 \tau_{\mathrm{C}}$ (see above). In this case (after the $90^{\circ}{ }^{1} \mathrm{H}$ pulse), the ligand and protein proton Zeeman reservoirs are aligned in an anti-parallel manner. Thus, the sign of the intermolecular NOESY cross peaks is inverted, whereas the intramolecular NOEs (within the ligand and the protein) are unchanged. Addition and subtraction give separate data sets for intra- and intermolecular NOEs, respectively. The two data sets were recorded in an interleaved manner and added and subtracted afterwards.


Figure 2. Simulation of the time evolution of the expectation values for transverse ${ }^{1} \mathrm{H}$ magnetization $\left\langle I_{x, y^{\prime}}\right\rangle$ during a ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ INEPT sequence. (A) Application of a hard ${ }^{13} \mathrm{C} 180^{\circ}$ pulse ( $5 \mu \mathrm{~s}$ ) optimized for a $J_{\mathrm{CH}}$ coupling of 120 Hz (applied at 6.35 ms ). (B) Adiabatic inversion pulse (sweep width: 60 kHz ; pulse duration: $2.8 \mathrm{~ms} ; \gamma \mathrm{B}_{1}$ : 5.0 kHz ; apodization of the first and last $30 \%$ using a sine function). The hard ${ }^{1} \mathrm{H}$ rf pulse ( $5 \mu \mathrm{~s}$ ) was applied in the middle of the INEPT step (after 5.26 ms ). The magnitude, $I_{x, y}=\left(I_{x}^{2}+I_{y}^{2}\right) 1 / 2$, of in-phase ${ }^{1} \mathrm{H}$ magnetization is shown. Simulations for three different $J_{\mathrm{CH}}$ values are shown: 120 Hz (solid); 140 Hz (dotted); 169 Hz (dashed-dotted).

## Results and discussion

The proposed sequence (Figure 3) was applied to the $\mathrm{B}_{12}$-binding subunit of glutamate mutase (MutS) from Clostridium tetanomorphum in complex with the nucleotide portion of the natural cofactor coenzyme $\mathrm{B}_{12}$. The solution structure of MutS has recently been published (Tollinger et al., 1998; Hoffmann et al., 2001) and revealed that the global fold of MutS in solution is similar to the crystal structures of other $\mathrm{B}_{12}$-binding domains (Drennan et al., 1994; Mancia et al., 1996; Reitzer et al., 1999). The MutS/B $\mathrm{B}_{12}$-nucleotide has recently been investigated using a ${ }^{15} \mathrm{~N}$-labeled MutS sample, the dissociation constant $K_{d}$ of the complex

Figure 3. Pulse scheme for simultaneous recording of intra- and intermolecular NOEs in protein-ligand complexes based on differential relaxation of longitudinal ${ }^{1} \mathrm{H}$ magnetization induced by selective ${ }^{1} \mathrm{H}$ inversion of protons attached to the ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled proteins. Experimental scheme to establish the difference mode ( $\Delta$ anti-parallel alignment of ${ }^{13} \mathrm{C}$, ${ }^{15} \mathrm{~N}$-protein and ligand ${ }^{1} \mathrm{H}$ ${ }_{13}$ gnetizations). The experiment for the generation of the sum mode ( $\Sigma$, parallel alignment of ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-protein and ligand ${ }^{1} \mathrm{H}$ magnetizations) is obtained by omitting the ${ }^{15} \mathrm{~N}$ and the adiabatic ${ }^{13} \mathrm{C}$ rf pulses in the boxed region (shown in grey). Narrow and wide pulses indicate $90^{\circ}$ and $180^{\circ}$ pulses, respectively, and, unless indicated otherwise, all pulses are applied along the $x$-axis.
Broadband ${ }^{13} \mathrm{C}$ inversion (during $\xi$ ) is achieved using a WURST adiabatic pulse (Kupce and Freeman, 1995) with a peak $\gamma \mathrm{B}_{1}=4.5 \mathrm{kHz}$ and with apodization of the first and last $30 \%$ using a ine function (Böhlen and Bodenhausen, 1993), a 60 kHz frequency sweep, and a duration of 2.8 ms . All rectangular ${ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}$ hard pulses (applied to ${ }^{13} \mathrm{C}_{\text {aliph }}$ aro) are applied with 28 and 20 kHz , respectively. ${ }^{15} \mathrm{~N}$ decoupling during $t_{1}$ is achieved with use of a WALTZ-16 field ( ${ }^{15} \mathrm{~N} \gamma \mathrm{~B}_{1}=840 \mathrm{~Hz}$ ) (Shaka et al., 1983). The ${ }^{13} \mathrm{C}$ ' shaped pulses have SEDUCE (McCoy and Müller, 1992) profiles ( $240 \mu \mathrm{~s}, 4.55 \mathrm{kHz}$ peak rf). The arrow indicates the position of the ${ }^{13} \mathrm{C}$ ' shaped pulse necessary for Bloch-Siegert compensation. The composite ${ }^{15} \mathrm{~N} 90^{\circ}$ rf pulses (before and after $t_{2}$ ) are applied to ensure a first-order phase correction of $180^{\circ}$ (Pascal et al., 1994). ${ }^{13} \mathrm{C}$ and ${ }^{15} \mathrm{~N}$ decoupling during acquisition were achieved with WURST (Kupce and Freeman, 1995) $\gamma \mathrm{B}_{1}=3.0 \mathrm{kHz}$ ) and GARP (Shaka et al., 1985) ( $\gamma \mathrm{B}_{1}=840 \mathrm{~Hz}$ ) decoupling scheme. The values for $\xi$ and $\tau_{\mathrm{a}}$, were set to 5.26 and 1.7 , respectively, and the NOESY mixing time $\tau_{\mathrm{m}}$ was set to 180 ms . Gradient levels were as follows: $\mathrm{g} 1=1.0 \mathrm{~ms}, 15 \mathrm{Gcm}^{-1} ; \mathrm{g} 2=1.0 \mathrm{~ms}, 8 \mathrm{Gcm}^{-1} ; \mathrm{g} 3=0.3 \mathrm{~ms}, 4 \mathrm{Gcm}^{-1} ; \mathrm{g} 4=1.0 \mathrm{~ms}, 15 \mathrm{Gcm}{ }^{-1} ; \mathrm{g} 5=1.0 \mathrm{~ms}, 20 \mathrm{Gcm}{ }^{-1} ; \mathrm{g} 6=1.0 \mathrm{~ms}, 8 \mathrm{Gcm}$;
$\mathrm{g} 7=4.0 \mathrm{~ms}, 19 \mathrm{Gcm}^{-1} ; \mathrm{g} 8=3.0 \mathrm{~ms},-18.8 \mathrm{Gcm}^{-1}$, and $\mathrm{g} 9=1.0 \mathrm{~ms}, 8 \mathrm{Gcm}^{-1}$. The phase cycling was $\phi 1=4(x), 4(-x) ; \phi 2=8(x), 8(-x) ; \phi 3=(x, y,-x,-y) ; \phi 4=4(x), 4(-x)$ and receiver was $(x,-y,-x, y), 2(-x, y, x,-y),(x,-y,-x, y)$. Quadrature detection in $\mathrm{F}_{1}$ and $\mathrm{F}_{2}$ is achieved by States-TPPI (Marion et al., 1989) of $\phi 1$ and $\phi 2$.


Figure 4. Illustration of the ${ }^{1} J_{\mathrm{CH}}$ compensation obtained with the pulse sequence element of Figure 1. 1D traces through F1 of the 3D ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ SIM-NOESY-HSQC (Figure 3). NOESY mixing time $\tau_{\mathrm{m}}$ was set to 180 ms . 1D traces showing the protein proton magnetizations (addition of the two data sets, intramolecular NOE data set) attached to carbon nuclei in the methyl (T126), $\mathrm{C} \alpha$ aliphatic (V133) and ${ }^{13} \mathrm{C}$ aromatic (Y64) region of the ${ }^{13} \mathrm{C}$ spectral range (thin line). Residual protein proton magnetization (subtraction of the two data sets, intermolecular NOE data set) (thick line). The ${ }^{13} \mathrm{C}$ chemical shifts of the selected carbons are as follows: T126(C $\gamma: 22.6 \mathrm{ppm})$, V133(C $\alpha: 64.2 \mathrm{ppm})$ and Y64(Cع:118.4ppm). Solid arrows indicate the diagonal peaks, dashed arrow indicate intermolecular NOEs, T126(H $\gamma$ )-H(C10/C11), see also Figure 6B, and NOEs to bulk water, respectively.
was determined to $5.6 \pm 0.7 \mathrm{mM}$ (Tollinger et al., 2001).

Experimental verification of the purging efficiency of the inversion element (Figure 1) is presented in Figure 4. 1D traces of selected carbons taken along the diagonal peaks of the intramolecular protein-protein (addition of the two data set) and corresponding traces from the intermolecular protein-ligand (subtraction of the two data set) NOESY map obtained with the sequence of Figure 3 are shown. Substantial suppression of protein diagonal peaks in the intermolecular

NOESY is achieved throughout the entire carbon spectrum $\left(\mathrm{CH}_{3}(\mathrm{~T} 126 \mathrm{H} \gamma): 4.8 \% ; \mathrm{C} \alpha(\mathrm{V} 133 \mathrm{H} \alpha): 8.4 \%\right.$ and $\left.\mathrm{C}_{\text {arom }}(\mathrm{Y} 64 \mathrm{H} \varepsilon): 4.2 \%\right)$. The experimental suppression levels are in qualitative agreement with the values calculated in Figure 2B. Most importantly, the trace of $\mathrm{T} 126 \mathrm{H} \gamma$ (Figure 4A), shows that both good suppression of protein ${ }^{1} \mathrm{H}$ signals and still intense intermolecular NOE cross-peaks are obtained with this experimental scheme. Although slightly better suppression levels might be achieved by employing more sophisticated (hyperbolic) frequency sweeps, for the sake of ease of implementation we suggest using linear sweeps.

The performance of the sequence was further validated by inspection of intramolecular NOEs between protons of ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled protein (Figure 5). For comparison we recorded a conventional $3 \mathrm{D}{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ -NOESY-HSQC (Figures 5a, 5c, and 5e) (Pascal et al., 1994). Representative expansions from the (F1,F3) NOE planes extracted at carbon frequencies typical for ${ }^{13} \mathrm{C} \alpha(\mathrm{F} 2=64.2$ and 56.5 ppm$)$ and ${ }^{13} \mathrm{C}$ methyl groups ( $\mathrm{F} 2=21.8 \mathrm{ppm}$ ) are shown (Figure 5). In both data sets essentially the same NOE connectivities were observed, indicating that the sequence can be used for recording intramolecular NOEs on the same protein sample under identical experimental conditions. The slight differential increase and decrease of NOE connectivities in panels e and f may be due to indirect water mediated dipolar relaxation pathways. In addition, the proposed sequence offers the possibility to discriminate between intermolecular NOEs to bound water molecules and intramolecular NOEs to H $\alpha$ 's with chemical shifts close to the water resonance frequency. It is noteworthy to compare our approach in terms of sensitivity with the previously proposed 4D $J$-resolved experiment (Melacini, 2000). In the $J$-resolved spectrum the signal intensity is reduced by a factor of 2 because the ${ }^{13} \mathrm{C}$-bound ${ }^{1} \mathrm{H}$ spins appear at $\pm{ }^{1} J_{\mathrm{CH}} / 2$ in the indirect frequency dimension and since the intramolecular NOESY spectrum is obtained by extracting the various NOE planes from the 4 D data set at $\mathrm{F} 2={ }^{1} J_{\mathrm{CH}} / 2$. The necessary spectral resolution (e.g. to resolve the ${ }^{13} \mathrm{C}-{ }^{1} \mathrm{H}$ doublet) along the J -dimension is achieved by adopting a constanttime (CT) version of the $J$-resolved experiment, which leads to an additional reduction due to ${ }^{1} \mathrm{H}$ transverse relaxation (Melacini, 2000). The overall length of the CT-J block in the $J$-resolved experiment was about 11.3 ms compared to $2 \xi=10.52 \mathrm{~ms}$ (Figure 3) in our case. A second important consideration is that the ${ }^{1} J_{\mathrm{CH}}$ scalar couplings are not uniform and thus an op-

$$
\delta(\mathrm{C} \alpha)=64.2 \mathrm{ppm}
$$




Figure 5. Identification of intramolecular protein NOEs. Representative plots from the ( $\mathrm{F}_{1}, \mathrm{~F}_{3}$ ) NOE plane extracted from the 3D experiment of Figure $3\left(\tau_{\mathrm{m}}=180 \mathrm{~ms}\right)$ after addition of the data set. ${ }^{13} \mathrm{C}$ spectral regions for methyl group $(21.8 \mathrm{ppm})$ and two selected $\mathrm{C} \alpha$ regions ( 56.5 and 64.2 ppm$)$ are shown. For comparison conventional ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ NOESY-HSQC spectra (Pascal et al., 1994) ( $\tau_{\mathrm{m}}=150 \mathrm{~ms}$ ) are also given. (a), (c) and (e) ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ NOESY-HSQC; (b), (d) and (f) intramolecular subset of ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ SIM-NOESY-HSQC (Figure 3). Acquisition and processing parameters are given in Materials and methods.
timal F2 selection cannot be simultaneously achieved for all ${ }^{13} \mathrm{C}$ bound protein protons.

An illustration of the observation of intermolecular NOEs is illustrated in Figure 6, where sections from the 3D SIM-NOESY-HSQC ( $\tau_{\mathrm{m}}=180 \mathrm{~ms}$ ) are shown. The data sets were recorded on an unlabeled $\mathrm{B}_{12}$-nucleotide-loop $/{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-MutS complex using the sequence of Figure 3 and by subtracting the two 3dimensional FIDs (sum and difference of protein and
ligand proton Zeeman reservoirs during $\tau_{\mathrm{m}}$ ). Figure 6 B shows three $\left({ }^{13} \mathrm{C},{ }^{1} \mathrm{H}\right)$ planes extracted at the ${ }^{1} \mathrm{H}$ chemical shifts (F1) of R5, Pr3 and C10/C11. the protons C10 and C11 are isochronous and cannot be distinguished. In Figure 6B correlations connecting the protons R 5 on the $\mathrm{B}_{12}$-nucleotide loop, with I22 21 , I2 $2 \gamma 2$ and V18 $\gamma 2$ on MutS are observed. Additionally, NOEs between Pr3 and I2281, V18 $\gamma 1$ and $\mathrm{V} 18 \gamma 2$, as well as between $\mathrm{C} 10 / \mathrm{C} 11$ and $\mathrm{T} 121 \gamma$,


Figure 6. (A) Molecular formula, numbering scheme and 1D spectrum of the nucleotide loop. (B) Selected $\left(\mathrm{F}_{2}, \mathrm{~F}_{3}\right)$ planes from the $3 \mathrm{D}{ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ SIM-NOESY-HSQC ( $\tau_{\mathrm{m}}=180 \mathrm{~ms}$ ) of unlabeled $\mathrm{B}_{12}$-nucleotide loop/ ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$-labeled MutS recorded with the sequence of Figure 3. Ligand signals are observed along $\mathrm{F}_{1}$, while protein signals are observed in $\mathrm{F}_{2}\left({ }^{13} \mathrm{C}\right)$ and $\mathrm{F}_{3}\left({ }^{1} \mathrm{H}\right) .\left({ }^{13} \mathrm{C},{ }^{1} \mathrm{H}\right)$ planes are taken at the ${ }^{1} \mathrm{H} \mathrm{F}_{1}$-frequency of the indicated protons ( $\mathrm{R} 5, \mathrm{Pr} 3$, and $\mathrm{C} 10 / \mathrm{C} 11$ ) of the nucleotide loop. In the middle panel ( Pr 3 ) of Figure 6B, the correlation peak marked with + has increased intensity in the adjacent slice and is due to a low molecular weight impurity. The intramolecular ligand cross-peak at natural ${ }^{13} \mathrm{C}$ abundance between the degenerate methyl groups C 10 and C 11 is indicated by ${ }^{* *}$ (lower panel, Figure $6 \mathrm{~B}, \mathrm{C} 10, \mathrm{C} 11$ ). Improved filtering of unwanted intramolecular protein signals in the intermolecular 3D NOESY-HSQC data set is demonstrated by $\mathrm{F}_{1}$ traces taken at the corresponding $\mathrm{F}_{2}\left({ }^{13} \mathrm{C}\right)$ and $\mathrm{F}_{3}\left({ }^{1} \mathrm{H}\right)$ frequencies of residues $\mathrm{I} 22 \gamma 2$, $\mathrm{I} 22 \delta 1$ and $\mathrm{T} 126 \gamma$ (indicated by arrows). Intermolecular and (remaining unwanted) intramolecular NOEs appear as positive and negative peaks, respectively, and are labeled either above (ligand) or below (protein) the horizontal traces.
$\mathrm{T} 125 \gamma$ and $\mathrm{T} 126 \gamma$ are found. The efficacy of the filtering scheme is additionally demonstrated in Figure 6B where traces for residues $\mathrm{I} 22 \gamma 2, \mathrm{I} 22 \delta 1$ and $\mathrm{T} 126 \gamma$ along the F 1 frequency dimension of the intermolecular 3D NOESY-HSQC data set are shown. Although only weak NOEs are detected in this low-affinity complex ( $K_{d}=5.6 \mathrm{mM}$, Tollinger et al., 2001), the suppression of intramolecular protein NOESY peaks is sufficient to ensure the unambiguous assignment of intermolecular NOEs. Additionally, note that remaining intramolecular diagonal peaks of the ${ }^{13} \mathrm{C},{ }^{15} \mathrm{~N}$ labeled protein are of opposite sign (with respect to the intermolecular cross peaks), whereas remaining intramolecular diagonal peaks of the ligand (at natural ${ }^{13} \mathrm{C}$ abundance) are of the same sign. All of the observed intermolecular NOEs could be unambiguously assigned based on the ${ }^{15} \mathrm{~N}$ solution structure of the $\mathrm{B}_{12}$-nucleotide-loop/MutS complex (Tollinger et al., 2001). The only exception, however, were weak NOEs observed between an (unassigned) aliphatic ${ }^{13} \mathrm{C}-{ }^{1} \mathrm{H}$ spin system ( $\left.{ }^{1} \mathrm{H}: 2.39 ;{ }^{13} \mathrm{C}: 33.6 \mathrm{ppm}\right)$ and methyl protons of I47, L74, I86, and L88. Interestingly, these side-chains form a hydrophobic cluster in the 3D structures of the apo-form of MutS (Hoffmann et al., 2001) and also in the MutS complex (Tollinger et al., 2001). The side-chains of these residues are adjacent to both a buried arginine side-chain (R75) and the thiol group of C78. We think that these NOEs are due to subtle changes in the intramolecular relaxation pathways of the protein protons mediated by exchangeable hydroxyl protons and/or internally bound water molecules which slightly alter the NOE cross peak intensities in the two experiments (sum or difference mode). We are currently investigating experimental possibilities to suppress these indirect relaxation pathways (e.g., suitable water inversion pulses in the middle of the NOESY mixing time applied according to the QUIET-NOESY technique, Zwahlen et al., 1994), as well as to possibly exploit this effect for hydration studies of proteins. It is important to note that the sequence of Figure 3 allows to record both intra- and intermolecular NOEs involving ${ }^{15} \mathrm{~N}$ bound amide protons. Specifically, the following intermolecular NOEs were observed: C10/C11- ${ }^{1} \mathrm{H}^{\mathrm{N}}$ (V90, G91, G92, Y117, T121, T126); C4-(V60, G91, G92); R5-(D24). Similar NOE connectivities were obtained on the ${ }^{15} \mathrm{~N}$-labeled sample (Tollinger et al., 2001).

The utility of the observed intermolecular NOEs obtained with the new sequence for structure determination of the MutS/B ${ }_{12}$-nucleotide loop complex is illustrated in Figure 7. In Figure 7A the X-ray
structure of the homologous protein glutamate mutase from clostridium cochlearium is shown together with the nucleotide moiety of the bound methylcobalamin (Reitzer et al., 1999). Figure 7B indicates the observed intermolecular NOEs obtained with the sequence of Figure 3. Sequence comparison between the $\mathrm{B}_{12}$-binding subunit of glutamate mutase from Clostridium cochlearium (GlmS) and MutS revealed that residues V125 and G126 in GlmS are replaced by T125 and T126. The structure used in Figure 7B was thus modelled accordingly. These NOEs provide critical restraints for establishing the correct binding mode of the nucleotide loop of the coenzyme to the $\mathrm{B}_{12}$-binding subunit of glutamate mutase, MutS. As we have discussed in some detail previously (Tollinger et al., 2001), ${ }^{15} \mathrm{~N}$ chemical shift perturbation and intermolecular NOEs (obtained from a ${ }^{15} \mathrm{~N}$-labeled protein sample) indicate a similar binding location of the $\mathrm{B}_{12}$-nucleotide loop in the MutS/B $\mathrm{B}_{12}$-nucleotide loop complex. For example, the proximity of the aromatic dimethylbenzimidazole in the $\mathrm{B}_{12}$-nucleotide loop to the central $\beta$-sheet ( $\beta 3$ and $\beta 4$, Figures 7A and 7B) was indicated by intermolecular NOEs between C10/C11, cf. Figure 6 A , and the amide protons ${ }^{1} \mathrm{H}^{\mathrm{N}}$ of S 62 ( $\beta 3$ ) and V90 and G91 ( $\beta 4$ ) (Tollinger et al., 2001). Additional intermolecular NOEs (between the protons C10/C11 and the ${ }^{1} \mathrm{H}^{\mathrm{N}}$ 's of I6, V60, Y117 and T126) further corroborated an analogous binding mode of the nucleotide in MutS, compared to the holoenzyme crystal structure of glutamate mutase (Reitzer et al., 1999). Interestingly, however, several NOEs could not be detected for ${ }^{15} \mathrm{~N}$-labeled MutS that would be expected based on the X-ray structure of glutamate mutase reconstituted with methylcobalamin (Reitzer et al., 1999) (e.g., H(C10/C11)- ${ }^{1} \mathrm{H}^{\mathrm{N}}$ (G9,C59); H(C4)${ }^{1} \mathrm{H}^{\mathrm{N}}(\mathrm{V} 60, \mathrm{~S} 61)$. Likewise, no intermolecular NOEs could be assigned for protons of the ribose ring or isopropanolamin unit of the $\mathrm{B}_{12}$-nucleotide. This indicated, together with the observed chemical shifts, a time averaged position of the $\mathrm{B}_{12}$-nucleotide somewhat more towards the entrance of the binding pocket (Tollinger et al., 2001). The enlarged number of unambiguously assigned intermolecular NOEs obtained with the sequence of Figure 1, in particular NOE connectivities between protons of the ribose R5 and the methyl group $\operatorname{Pr} 3$ of the isopropanolamin unit and protons of MutS allow now a better characterization of the binding site.

It is also important to note, that due to the high molar concentration ( $\approx 24 \mathrm{mM}$ ) of the unlabeled $\mathrm{B}_{12^{-}}$ nucleotide loop, intramolecular $\mathrm{B}_{12}$-nucleotide NOE


Figure 7. Ribbon diagramm showing the $\mathrm{B}_{12}$-nucleotide binding cleft of MutS. (A) Structural homology model for the $\mathrm{B}_{12}$-nucleotide binding site in MutS (based on the crystal structure of glutamate mutase from Clostridium cochlearium, Reitzer et al., 1999). The location of the $\mathrm{B}_{12}$-nucleotide binding site in MutS was obtained from the closely related complex between glutamate mutase from Clostridium cochlearium and methylcobalamin (see text) (Reitzer et al., 1999) and independently confirmed in solution by ${ }^{15} \mathrm{~N}$ chemical shift perturbation (Tollinger et al., 2001). Secondary structure elements are numbered. Only the nucleotide moiety of bound methylcobalamin is shown. (B) Schematic diagram showing the helices $\alpha 1$ and $\alpha 5$ and $\beta$-strands $\beta 1, \beta 3$ and $\beta 4$ of MutS and the intermolecular MutS/ $\mathrm{B}_{12}$-nucleotide NOEs observed with the sequence of Figure 3 (dashed lines). The figure was prepared using the program MOLMOL (Koradi et al., 1996).
cross peaks at natural abundance become visible in the intermolecular NOE data set (see Figure 6B). For example, in the lower panel of Figure 6B, the cross peak denoted by ${ }^{* *}$ is due to an intramolecular NOE between the degenerate methyl groups C10 and C11 of the dimethylbenzimidazole moiety. However, these intramolecular cross peaks can also be used for structural refinement purposes (e.g., to better characterize the geometry of the bound ligand). Specifically, we observed the following intramolecular dipolar connectivities: C7-R1, C7-C11, R5-R4, Pr2-Pr3, and $\operatorname{Pr} 1-\operatorname{Pr} 3$, which will be beneficial for the refinement of the protein bound conformation (e.g., sugar pucker, glycosidic angle) of the $\mathrm{B}_{12}$-nucleotide loop.

## Conclusions

A novel method for simultaneous recording of interand intramolecular NOEs in molecular complexes consisting of ${ }^{15} \mathrm{~N},{ }^{13} \mathrm{C}$-labeled proteins and unlabeled ligands has been presented. The method relies on
a selective inversion of protein protons (attached to ${ }^{13} \mathrm{C}$ and ${ }^{15} \mathrm{~N}$ ) by applying an adiabatic ${ }^{13} \mathrm{C}$ frequency sweep optimized for the empirically observed relationship between the heteronuclear one-bond scalar coupling constant ${ }^{1} J_{\mathrm{CH}}$ and the ${ }^{13} \mathrm{C}$ chemical shift. Two data sets are recorded in which the relative signs of ligand and protein Zeeman magnetizations during the NOESY mixing period are inverted. Addition or subtraction of the two data sets results in pure intramolecular or intermolecular NOESY maps, respectively. The particular merits of the proposed filtering method is the significant ${ }^{1} J_{\mathrm{CH}}$-bandwidth and a substantial time saving by allowing the simultaneous measurement of intra- and intermolecular NOEs. Data obtained on the low-affinity complex between the $\mathrm{B}_{12}$-binding subunit of glutamate mutase from Clostridium tetanomorphum and the $\mathrm{B}_{12}$-nucleotide loop also suggest possible applications to studies of macromolecular complexes in the context of SAR by NMR (Shuker et al., 1996).

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